

Titre: Mutual Inductance and Magnetic Force Calculations for Bitter Disk Coil (Pancake) with Nonlinear Radial Current and Filamentary Circular Coil with Azimuthal Current
Title:

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Authors:

Date: 2016

Type: Article de revue / Article

Référence: Babic, S., & Akyel, C. (2016). Mutual Inductance and Magnetic Force Calculations for Bitter Disk Coil (Pancake) with Nonlinear Radial Current and Filamentary Circular Coil with Azimuthal Current. *Advances in Electrical Engineering*, 2016, 1-6. <https://doi.org/10.1155/2016/3654021>
Citation:

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 **Document publié chez l'éditeur officiel**
Document issued by the official publisher

Titre de la revue: Advances in Electrical Engineering (vol. 2016)
Journal Title:

Maison d'édition: Hindawi Publishing Corporation
Publisher:

URL officiel: <https://doi.org/10.1155/2016/3654021>
Official URL:

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Research Article

Mutual Inductance and Magnetic Force Calculations for Bitter Disk Coil (Pancake) with Nonlinear Radial Current and Filamentary Circular Coil with Azimuthal Current

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Received 12 June 2016; Revised 8 August 2016; Accepted 22 August 2016

Academic Editor: Gorazd Stumberger

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Bitter coils are electromagnets used for the generation of extremely strong magnetic fields superior to 30 T. In this paper we calculate the mutual inductance and the magnetic force between Bitter disk (pancake) coil with the nonlinear radial current and the circular filamentary coil with the azimuthal current. The close form expressed over complete elliptic integrals of the first and second kind as well as Heuman's Lambda function is obtained for this configuration either for the mutual inductance or for the magnetic force. The results of this method are compared with those obtained by the improved modified filament method for the presented configuration. All results are in an excellent agreement.

1. Introduction

In the literature and scientific papers the calculation of the mutual inductance and the magnetic force between ordinary circular coils with the uniform current densities have been given by different methods either analytical or numerical [1–15]. However, there are not a lot references about the mutual inductance and the magnetic force calculation between circular coils in which the current densities are not uniform. The coils with this characteristic are Bitter coils or Bitter solenoids. Bitter coils are used in high magnetic field applications and they differ from the ordinary coils in that they have an inverse radial distribution of current [16, 17]. The interesting point is the calculation of the mutual inductance and the magnetic force between Bitter coils or between one Bitter coil and an ordinary coil. Generally, the calculation of the magnetic force between circular coils is closely related to the calculation of their mutual inductance. Since their mutual energy is equal to the product of their mutual inductance and the currents in the coils, the component of the magnetic force of attraction or repulsion in any direction is equal to the product of the currents multiplied by the differential

coefficient of the mutual inductance taken with respect to that coordinate. The mutual inductance and the magnetic force are obtained over multiple integrals with the different kernel functions. Thus, in the calculation of the mutual inductance and the corresponding magnetic force between two coils we need to integrate their kernel functions depending on coil configurations. Hopefully, these kernel functions are Green functions $r^{-0.5}$ and $r^{-1.5}$, where the “ r ” is the distance between two coils. This integral approach, much easier than the differential approach, leads to relatively simple expressions which are incorporated in these two physical quantities. In this paper we calculate the mutual inductance between the Bitter disk coil with nonlinear radial current and a circular filamentary coil with the azimuthal current (ordinary coil). Coils are coaxial and in air. Either the mutual inductance or the magnetic force is obtained in closed form expressed over complete elliptic integrals of the first and second kind as well as over Heuman's Lambda function [18, 19]. The results of these calculations will be compared to those obtained by the improved filament method for concerned configuration. The results obtained by these two methods are in an excellent agreement.

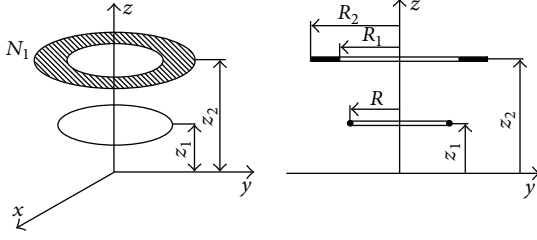


FIGURE 1: The thin Bitter disk coil (pancake) and the filamentary circular coil.

2. Basic Expressions

The mutual inductance and magnetic force between a disk coil (pancake) with the uniform azimuthal current density $J_1 = N_1 I_1 / (R_2 - R_1)$ and a filamentary coil with the azimuthal current I_2 (see Figure 1) can be calculated, respectively, by [1, 2]

$$M = \frac{\mu_0 N_1 R}{(R_2 - R_1)} \int_0^\pi \int_{R_1}^{R_2} \frac{r_I \cos \theta dr_I d\theta}{r_0}, \quad (1)$$

$$F = -\frac{\mu_0 N_1 I_1 I_2 R (z_2 - z_1)}{(R_2 - R_1)} \int_0^\pi \int_{R_1}^{R_2} \frac{r_I \cos \theta dr_I d\theta}{r_0^3}, \quad (2)$$

where

$$r_0 = \sqrt{(z_2 - z_1)^2 + r_I^2 + R^2 - 2r_I R \cos \theta} \quad (3)$$

and N_1 is the number of turns of the disk coil (pancake).

$\mu_0 = 4\pi \times 10^{-7}$ H/m is the permeability of free space (vacuum), and r, θ, z are the cylindrical coordinates. Both coils are in air or in a nonmagnetic and nonconducting environment.

In [2] the magnetic force is obtained in the close form expressed over the complete elliptic integral of the first kind and Heuman's Lambda function and the magnetic force in the semianalytical form expressed over complete elliptic integral of the first and second kind, Heuman's Lambda function plus one simple integral which has to be evaluated by some of the numerical integrations [1].

Let us suppose the disk coil be the Bitter coil in which the current density is not uniform [16, 17] and given by

$$J_1 = \frac{N_1 I_1}{r_I \ln(R_2/R_1)}. \quad (4)$$

From (1), (2), and (4) we obtain the expressions of the mutual inductance and the magnetic force between the Bitter coil and the filamentary coil as follows:

$$M_B = \frac{\mu_0 N_1 R}{\ln(R_2/R_1)} \int_0^\pi \int_{R_1}^{R_2} \frac{\cos \theta dr_I d\theta}{r_0}, \quad (5)$$

$$F_B = -\frac{\mu_0 N_1 I_1 I_2 R (z_2 - z_1)}{\ln(R_2/R_1)} \int_0^\pi \int_{R_1}^{R_2} \frac{\cos \theta dr_I d\theta}{r_0^3}, \quad (6)$$

where

$$r_0 = \sqrt{(z_2 - z_1)^2 + r_I^2 + R^2 - 2r_I R \cos \theta}. \quad (7)$$

3. Calculation Method

Integrating in (5) and (6) over r_I and θ (substituting $\theta = \pi - \beta$) the mutual inductance and the magnetic force between the Bitter disk and the filamentary circular coil can be expressed, respectively, in an analytical form as follow (see Appendices A and B):

$$M_B = \frac{\mu_0 N_1}{2 \ln(R_2/R_1)} \sum_{n=1}^{n=2} (-1)^{n-1} T_n, \quad (8)$$

$$F_B = \frac{\mu_0 N_1 I_1 I_2}{\ln(R_2/R_1)} \sum_{n=1}^{n=2} (-1)^{n-1} S_n, \quad (9)$$

where

$$T_n = \frac{k_n}{\sqrt{RR_n}} \left[2R \sqrt{R^2 + z_Q^2} - R_n^2 - R^2 - z_Q^2 \right] K(k_n) + \frac{k_n}{\sqrt{RR_n}} \left[(R + R_n)^2 + z_Q^2 \right] E(k_n) - \pi |z_Q| V_n,$$

$$S_n = \frac{k_n z_Q}{\sqrt{RR_n}} \frac{R}{\sqrt{R^2 + z_Q^2} + R} K(k_n) - \frac{\pi}{2} \operatorname{sgn}(z_Q) V_n,$$

$$V_n = 1 - \Lambda_0(\theta_{1n}, k_n)$$

$$+ \operatorname{sgn}\left(\sqrt{R^2 + z_Q^2} - R_n\right) \left[1 - \Lambda_0(\theta_{2n}, k_n) \right],$$

$$z_Q = z_2 - z_1, \quad (10)$$

$$k_n^2 = \frac{4RR_n}{(R + R_n)^2 + z_Q^2} \leq 1,$$

$$m = \frac{2R}{\sqrt{R + z_Q^2} + R} \leq 1,$$

$$\theta_{1n} = \arcsin \frac{|z_Q|}{\sqrt{R^2 + z_Q^2} + R},$$

$$\theta_{2n} = \arcsin \sqrt{\frac{1-m}{1-k_n^2}}, \quad k_n^2 \leq m.$$

Thus, all expressions are obtained in the closed form expressed over complete elliptic integrals of the first and second kind as well as Heuman's Lambda function [18, 19].

3.1. Special Cases and Singularities

Mutual Inductance. In the mutual inductance calculation special cases and singular cases appear for $z_Q = 0$ and $k_n^2 \neq 1$ or $z_Q = 0$ and $k_n^2 = 1$.

3.1.1. $z_Q = 0$ and $k_n^2 \neq 1$. This case is the regular case covered by (8). In this case the term T_n can be given in the following simplified form:

$$T_n = \frac{k_{n0} (R + R_n)}{\sqrt{RR_n}} [(R - R_n) K(k_{n0}) + (R + R_n) E(k_{n0})], \quad (11)$$

$$k_{n0}^2 = \frac{4RR_n}{(R + R_n)^2} \leq 1. \quad (12)$$

3.1.2. $z_Q = 0$ and $k_n^2 = 1$. This case is the singular case ($R = R_1$ or $R = R_2$) for which the term T_n is equal either $4R_1$ or $4R_2$.

Magnetic Force. In the magnetic force calculation for $z_Q = 0$ its value is zero because the coils are in the same plane.

Notation Above. Λ_0 corresponds to Heuman's Lambda function as defined in [18, 19] and "sign" function returns integer indicating the sign of a number. Functions K and E correspond to complete elliptic integrals of the first and second kind [18, 19].

4. Modified Filament Method

In [17] the mutual inductance is calculated between the Bitter coil and the superconducting coil with uniform current density by using the filament method. Here we give the modified formulas for the mutual inductance and the magnetic force for the treated configuration (see Figure 2) using the filament method.

Applying some modification in the mutual inductance calculation given in [17] we deduced the mutual inductance and the magnetic force between the Bitter coil and the filamentary coil as follows:

$$M_{BF} = \frac{N_1 (R_2 - R_1)}{(2n + 1) \ln(R_2/R_1)} \sum_{l=-n}^{l=n} \frac{M(l)}{r_I(l)}, \quad (13)$$

$$F_{BF} = \frac{N_1 I_1 I_2 (R_2 - R_1)}{(2n + 1) \ln(R_2/R_1)} \sum_{l=-n}^{l=n} \frac{F(l)}{r_I(l)},$$

where

$$M(l) = \frac{2\mu_0 \sqrt{r_I(l) R}}{k(l)} \left[\left(1 - \frac{k^2(l)}{2} \right) K(k(l)) - E(k(l)) \right],$$

$$F(l) = -\frac{\mu_0 I_1 I_2 c k(l)}{4\sqrt{r_I(l) R}} \left[\frac{2 - k^2(l)}{1 - k^2(l)} E(k(l)) - 2K(k(l)) \right],$$

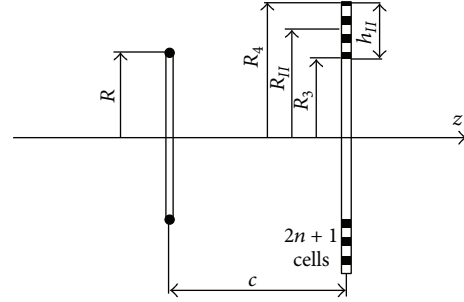


FIGURE 2: Configuration of mesh matrix: Filamentary coil-thin disk coil (pancake).

$$r_I(l) = \frac{R_2 + R_1}{2} + \frac{R_2 - R_1}{2n + 1} l \quad (l = -n, \dots, 0, \dots, n),$$

$$c = z_2 - z_1 = z_Q,$$

$$k^2(l) = \frac{4r_I(l) R}{(r_I(l) + R)^2 + c^2}. \quad (14)$$

Expressions (13) will be used to confirm the validity of analytical formulas (8) and (9).

5. Examples

To verify the validity of the new formulas, we apply them to the following set of examples.

Example 1. Calculate the mutual inductance and the magnetic force between the Bitter disk coil and a filamentary circular coil with the following dimensions and the number of turns. Currents in coils are unit.

Filamentary Circular Coil. $R = 20$ mm, $z_1 = 0$ mm.

Bitter Disk Coil. $R_1 = 40$ mm, $R_2 = 60$ mm, $z_2 = 50$ mm, $N_1 = 100$.

Applying (8) and (9) the mutual inductance and the magnetic force are, respectively,

$$M_B = 521.47220251309271 \text{ nH},$$

$$F_B = -16.016662833788816 \mu\text{N}. \quad (15)$$

The computational time is about 0.037677 seconds.

By using the modified filament, (13) and (14), we obtain

$$M_{BF} = 521.4722025131178 \text{ nH},$$

$$F_{BF} = -16.01666283378895 \mu\text{N}. \quad (16)$$

The number of subdivision was $n = 1000000$ and the computational time about 0.362069 seconds.

Example 2. Calculate the mutual inductance between a disk and a filamentary circular coil with the following dimensions and the number of turns.

Filamentary Circular Coil. $R = 20$ mm, $z_1 = 0$ mm.

Disk Coil. $R_1 = 40$ mm, $R_2 = 60$ mm, $z_2 = 0$ mm, $N_1 = 100$.

Applying (8) and (9) the mutual inductance and the magnetic force are, respectively,

$$\begin{aligned} M_B &= \mathbf{1.7442876932070914} \mu\text{H}, \\ F_B &= 0 \text{ N}. \end{aligned} \quad (17)$$

The computational time is about 0.020657 seconds.

By using the modified filament, (13) and (14), we obtain

$$\begin{aligned} M_{\text{BF}} &= \mathbf{1.744287693207058} \mu\text{H}, \\ F_{\text{BF}} &= 0 \text{ N}. \end{aligned} \quad (18)$$

The number of subdivision was $n = 20000$ and the computational time about 0.030564 seconds.

Example 3. Calculate the mutual inductance between a disk and a filamentary circular coil with the following dimensions and the number of turns.

Filamentary Circular Coil. $R = 40$ mm, $z_1 = 0$ mm.

Disk Coil. $R_1 = 40$ mm, $R_2 = 60$ mm, $z_2 = 0$ mm, $N_1 = 100$.

This case is the singular case because $R = R_1$.

Applying (8), (9), and (11) the mutual inductance and the magnetic force are, respectively,

$$\begin{aligned} M_B &= \mathbf{10.931677144024915} \mu\text{H}, \\ F_B &= 0 \text{ N}. \end{aligned} \quad (19)$$

The computational time is about 0.020223 seconds.

By using the modified filament, (13) and (14), we obtain

$$\begin{aligned} M_{\text{BF}} &= \mathbf{10.93167597572536} \mu\text{H}, \\ F_{\text{BF}} &= 0 \text{ N}. \end{aligned} \quad (20)$$

The number of subdivision was $n = 1000000$ and the computational time about 0.513017 seconds.

Example 4. Calculate the mutual inductance between a disk and a filamentary circular coil with the following dimensions and the number of turns.

Filamentary Circular Coil. $R = 60$ mm, $z_1 = 0$ mm.

Disk Coil. $R_1 = 40$ mm, $R_2 = 60$ mm, $z_2 = 0$ mm, $N_1 = 100$.

This case is also the singular case for which $R = R_2$.

Applying (8), (9), and (11) the mutual inductance and the magnetic force are, respectively,

$$\begin{aligned} M_B &= \mathbf{14.062056002507577} \mu\text{H}, \\ F_B &= 0 \text{ N}. \end{aligned} \quad (21)$$

The computational time is about 0.03929 seconds.

By using the modified filament, (13) and (14), we obtain

$$\begin{aligned} M_{\text{BF}} &= \mathbf{14.0620549882661} \mu\text{H}, \\ F_{\text{BF}} &= 0 \text{ N}. \end{aligned} \quad (22)$$

The number of subdivision was $n = 1000000$ and the computational time about 0.487595 seconds.

In comparative calculations the significant figures which agree are bolded.

Thus, with these examples we confirmed the validity of the present analytical method for calculating the mutual inductance and the magnetic force between the Bitter disk with the nonlinear radial current distribution and an ordinary filamentary circular coil. All calculations were executed in Matlab programming.

6. Conclusion

The new accurate mutual inductance and magnetic force formulas for the system of the Bitter disk with the nonlinear radial current and the filamentary circular with the azimuthal in air are derived and presented in this paper. All expressions for either the mutual inductance or the magnetic force are obtained in the close form expressed over complete elliptic integrals of the first and second kind and Heuman's Lambda function. Also we gave in this paper the improved formulas for the mutual inductance and the magnetic force between treated coils by using the filament method. All singular cases are obtained in a closed form. The presented method can be used to calculate these important electrical quantities such as the mutual inductance and the magnetic force for coils combinations comprising Bitter coils of rectangular cross section and the circular coils with uniform current densities.

Appendix

A. Mutual Inductance Integrals

The first integral in (5) is I_{10} which can be solved analytically [19]:

$$\begin{aligned} I_{10} &= \int_{R_1}^{R_2} \frac{dr_I}{\sqrt{r_I^2 + 2r_I R \cos 2\beta + R^2 + z_Q^2}} \\ &= \text{arsh} \frac{r_I + R \cos 2\beta}{\sqrt{R^2 \sin^2 2\beta + z_Q^2}} \Big|_{R_1}^{R_2}. \end{aligned} \quad (\text{A.1})$$

The last integral in (5) is

$$I_{20} = \int_0^{\pi/2} \cos 2\beta \cdot \text{arsh} \frac{R_n + R \cos 2\beta}{\sqrt{R^2 \sin^2 2\beta + z_Q^2}} d\beta \quad (\text{A.2})$$

which can be obtained in the following form:

$$I_{20} = \frac{k_n}{2} \sqrt{\frac{R_n}{R}} \int_0^{\pi/2} \frac{\cos^3 2\beta + \left(\frac{R^2 + z_Q^2}{RR_n}\right) \cos^2 2\beta - \cos 2\beta - \left(\frac{R^2 + z_Q^2}{RR_n}\right) \Delta}{\left(\cos^2 2\beta - \left(\frac{R^2 + z_Q^2}{R^2}\right) \Delta\right) \Delta} d\beta = \frac{k_n}{2} \sqrt{\frac{R_n}{R}} I_0, \quad (\text{A.3})$$

where

$$z_Q = z_2 - z_1,$$

$$k_n^2 = \frac{4RR_n}{(R + R_n)^2 + z_Q^2} \leq 1,$$

$$\Delta = \sqrt{1 - k_n^2 \sin^2 2\beta}.$$

(A.4)

The integral I_0 can be obtained in its final form expressed over complete elliptic integrals of the first and second kind $K(k)$, $E(k)$ as well as Heuman's Lambda function $\Lambda_0(\epsilon, k)$ [18, 19]:

$$\begin{aligned} I_0 &= \int_0^{\pi/2} \frac{\cos^3 2\beta + \left(\frac{R^2 + z_Q^2}{RR_n}\right) \cos^2 2\beta - \cos 2\beta - \left(\frac{R^2 + z_Q^2}{RR_n}\right) \Delta}{\left(\cos^2 2\beta - \left(\frac{R^2 + z_Q^2}{R^2}\right) \Delta\right) \Delta} d\beta = \int_0^{\pi/2} \frac{\cos 2\beta + \left(\frac{R^2 + z_Q^2}{RR_n}\right) \Delta}{\Delta} d\beta \\ &+ \frac{z_Q^2}{R^2} \int_0^{\pi/2} \frac{\cos 2\beta + \left(\frac{R^2 + z_Q^2}{RR_n}\right) \Delta}{\left(\cos^2 2\beta - \left(\frac{R^2 + z_Q^2}{R^2}\right) \Delta\right) \Delta} d\beta = \frac{2R\sqrt{R^2 + z_Q^2} - R_n^2 - R^2 - z_Q^2}{2RR_n} K(k_n) + \frac{(R + R_n)^2 + z_Q^2}{2RR_n} E(k_n) \\ &- \frac{\pi |z_Q|}{2k_n \sqrt{RR_n}} V_n, \end{aligned} \quad (\text{A.5})$$

where

$$\begin{aligned} V_n &= 1 - \Lambda_0(\theta_{1n}, k_n) \\ &+ \operatorname{sgn}\left(\sqrt{R^2 + z_Q^2} - R_n\right) [1 - \Lambda_0(\theta_{2n}, k_n)]. \end{aligned} \quad (\text{A.6})$$

Replacing I_{20} in expression (5) we obtain (8).

B. Magnetic Force Integrals

The first integral in (6) is L_{10} which has the analytic solution [19]:

$$\begin{aligned} L_{10} &= \int_{R_1}^{R_2} \frac{dr_I}{\left(r_I^2 + 2r_I R \cos 2\beta + R^2 + z_Q^2\right)^{3/2}} \\ &= \frac{r_I + R \cos 2\beta}{\left(R^2 \sin^2 2\beta + z_Q^2\right) \sqrt{r_I^2 + 2r_I R \cos 2\beta + R^2 + z_Q^2}} \Bigg|_{R_1}^{R_2}. \end{aligned} \quad (\text{B.1})$$

The last integral in (6) is

$$\begin{aligned} L_{20} &= 2Rz_Q \int_0^{\pi/2} \frac{\cos 2\beta (R_n + R \cos 2\beta) d\beta}{\left(R^2 \sin^2 2\beta + z_Q^2\right) \sqrt{r_I^2 + 2r_I R \cos 2\beta + R^2 + z_Q^2}} \end{aligned} \quad (\text{B.2})$$

which can be given in the following form:

$$\begin{aligned} L_{20} &= -\frac{k_n z_Q}{\sqrt{R_n R}} \left[\int_0^{\pi/2} \frac{d\beta}{\Delta} \right. \\ &+ \left. \int_0^{\pi/2} \frac{\cos 2\beta + \left(\frac{R^2 + z_Q^2}{RR_n}\right) \Delta}{\left(\cos^2 2\beta - \left(\frac{R^2 + z_Q^2}{R^2}\right) \Delta\right) \Delta} d\beta \right] \\ &= -\frac{k_n z_Q}{\sqrt{R_n R}} \left[\int_0^{\pi/2} \frac{d\beta}{\Delta} + \frac{R_n}{R} I_0 \right] \\ &= -\frac{k_n z_Q}{\sqrt{R_n R}} \left[K(k_n) + \frac{R_n}{R} I_0 \right]. \end{aligned} \quad (\text{B.3})$$

Integral I_0 appears in Appendix A so that replacing L_{20} in (6) we obtain (9).

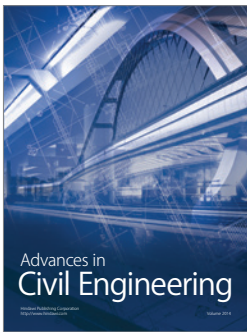
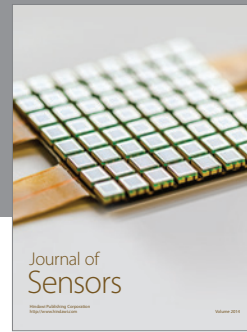
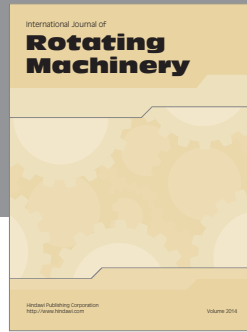
Competing Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

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